

Article

# DC Self-Field Critical Current in Superconductor/Dirac-Cone Material/Superconductor Junctions

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Received: 25 September 2019; Accepted: 29 October 2019; Published: 1 November 2019



**Abstract:** Recently, several research groups have reported on anomalous enhancement of the self-field critical currents,  $I_c(sf, T)$ , at low temperatures in superconductor/Dirac-cone material/superconductor (S/DCM/S) junctions. Some papers attributed the enhancement to the low-energy Andreev bound states arising from winding of the electronic wave function around DCM. In this paper,  $I_c(sf, T)$  in S/DCM/S junctions have been analyzed by two approaches: modified Ambegaokar-Baratoff and ballistic Titov-Beenakker models. It is shown that the ballistic model, which is traditionally considered to be a basic model to describe  $I_c(sf, T)$  in S/DCM/S junctions, is an inadequate tool to analyze experimental data from these type of junctions, while Ambegaokar-Baratoff model, which is generally considered to be a model for  $I_c(sf, T)$  in superconductor/insulator/superconductor junctions, provides good experimental data description. Thus, there is a need to develop a new model for self-field critical currents in S/DCM/S systems.

**Keywords:** the self-field critical current; induced superconductivity in Dirac-cone materials; single layer graphene; multiple-band superconductivity

## 1. Introduction

Intrinsic superconductors [1] of rectangular cross-section (with width  $2a$  and thickness  $2b$ ) exhibit non-dissipative temperature dependent transport self-field critical current,  $I_c(sf, T)$  (i.e., when no external magnetic field applies), which is given by the following universal equation [2–4]:

$$I_c(sf, T) = \frac{\phi_0}{\pi \mu_0} \left[ \frac{\ln(1 + \sqrt{2} \kappa_c(T))}{\lambda_{ab}^3(T)} \left( \frac{\lambda_c(T)}{b} \tanh\left(\frac{b}{\lambda_c(T)}\right) \right) + \frac{\ln(1 + \sqrt{2} \gamma(T) \kappa_c(T))}{\sqrt{\gamma(T)} \lambda_{ab}^3(T)} \left( \frac{\lambda_{ab}(T)}{a} \tanh\left(\frac{a}{\lambda_{ab}(T)}\right) \right) \right] \cdot (a \cdot b), \quad (1)$$

where  $T$  is sample temperature,  $\phi_0 = 2.067 \times 10^{-15}$  Wb is the magnetic flux quantum,  $\mu_0 = 4 \cdot \pi \times 10^{-7}$  H/m is the magnetic permeability of free space,  $\lambda_{ab}(T)$  and  $\lambda_c(T)$  are the in-plane and out-of-plane London penetration depths respectively,  $\kappa_c(T) = \lambda_{ab}(T) / \xi_{ab}(T)$ ,  $\xi_{ab}(T)$  is the in-plane coherence length, and  $\gamma(T) = \lambda_c(T) / \lambda_{ab}(T)$  is the electron mass anisotropy. It has been shown in previous research that Equation (1) quantitatively and accurately describes  $I_c(sf, T)$  in more than 100 superconductors, ranging from elemental Zn with  $T_c = 0.65$  K to highly-compressed H<sub>3</sub>S with  $T_c \gtrsim 200$  K [2–4], and samples dimensions from several Å to about 1 mm [5].

All intrinsic superconductors [1] can induce a superconducting state in non-superconducting materials by the Holm-Meissner effect [6]. However, a universal equation for non-dissipative self-field critical transport current,  $I_c(sf, T)$ , in superconductor/non-superconductor/superconductor junctions is still unknown. Ambegaokar and Baratoff (AB) [7,8] were the first who proposed an equation for  $I_c(sf, T)$  in superconductor/insulator/superconductor (S/I/S) systems [9]. Later, Kulik and

Omel'yanchuk (KO) [10–12] proposed two models for different types of superconductor/normal conductor/superconductor junctions (which are known as KO-1 [10] and KO-2 [11]).

In general, superconductor/normal metal/superconductor (S/N/S) junctions are classified by the comparison of the device length ( $L$ ) to two characteristic length scales of the junction, which are the mean free path of the charge carriers,  $l_e$ , and the superconducting correlation length,  $\xi_s$ . These length scales classify whether the junction is in short ( $L \ll \xi_s$ ) or long (i.e.,  $L \gg \xi_s$ ) regime and ballistic ( $L \ll l_e$ ) or diffusive ( $L \gg l_e$ ) limit, respectively.

For about one decade, the KO-1 model was considered to be the primary model to describe  $I_c(sf, T)$  in superconductor/graphene/superconductor (S/G/S) junctions (a detailed review of different models for  $I_c(sf, T)$  in S/G/S junctions was given by Lee and Lee [13]). However, recent technological progress in fabricating high-quality S/G/S junctions demonstrates a large difference between the KO-1 model and experimental  $I_c(sf, T)$  data [14]. A detailed discussion of all models, including a model by Takane and Imura [15], which was proposed to describe  $I_c(sf, T)$  in superconductor/Dirac-cone material/superconductor (S/DCM/S) junctions, is given by Lee and Lee [13].

It should be noted that a universal quantitatively accurate equation for critical currents at the applied magnetic field,  $B$ , is unknown to date for intrinsic superconductors [16–20] and for Josephson junctions [13,21,22]. However, the discussion of these important problems, as well as the discussion of interface superconductivity [23–25] and generic case of two-dimensional (2D) superconductivity [26–50], is beyond the scope of this paper.

The primary task for this work is to show that  $I_c(sf, T)$ , in a variety of S/DCM/S junctions in the ballistic regime, cannot be described by the KO-based model. To prove this, experimental  $I_c(sf, T)$  datasets in S/DCM/S junctions were analyzed by two models: the modified Ambegaokar-Baratoff model [51,52] and ballistic Titov-Beenakker model [53].

It needs to be noted that some S/DCM/S junctions show the  $I_c(sf, T)$  enhancement at a reduced temperature of  $T \leq 0.25 \cdot T_c$ . For instance, the enhancement in atomically-thin MoRe/single layer graphene (SLG)/MoRe junction was first reported by Calado et al. [54]. Raw experimental  $I_c(sf, T)$  data reported by Borzenets et al. [55] in nominally the same MoRe/SLG/MoRe junctions also shows the enhancement at  $T \leq 0.25 \cdot T_c$ . Based on this, the  $I_c(sf, T)$  enhancement at low reduced temperatures in Nb/BiSbTeSe<sub>2</sub>-nanoribbon/Nb reported by Kayyalha et al. [56] cannot be considered as a unique property of superconductor/topological insulator/superconductor (S/TI/S) junctions, but is rather the demonstration of a general feature of S/DCM/S devices and atomically thin superconducting systems. Additionally, it is important to mention that Kurter et al. [57] were the first who reported  $I_c(sf, T)$  enhancement in S/TI-nanoribbon/S junction at reduced temperature of  $T \leq 0.25 \cdot T_c$ .

As a result of the performed  $I_c(sf, T)$  analysis in this paper, it is shown that a new model is needed to describe dissipation-free transport currents in S/DCM/S junctions.

## 2. Models Description

The amplitude of dissipation-free transport current,  $I_c(sf, T)$ , in S/I/S junction was first given by Ambegaokar and Baratoff (AB) [7,8]:

$$I_c(sf, T) = \frac{\pi \cdot \Delta(T)}{2 \cdot e \cdot R_n} \cdot \tanh\left(\frac{\Delta(T)}{2 \cdot k_B \cdot T}\right), \quad (2)$$

where  $\Delta(T)$  is the temperature-dependent superconducting gap,  $e$  is the electron charge,  $R_n$  is the normal-state tunneling resistance in the junction, and  $k_B$  is the Boltzmann constant. In one research [51], it was proposed to substitute  $\Delta(T)$  in Equation (2) by the analytical expression given by Gross et al. [58]:

$$\Delta(T) = \Delta(0) \cdot \tanh\left(\frac{\pi \cdot k_B \cdot T_c}{\Delta(0)} \cdot \sqrt{\eta \cdot \left(\frac{\Delta C}{C}\right) \cdot \left(\frac{T_c}{T} - 1\right)}\right), \quad (3)$$

where  $\Delta(0)$  is the ground-state amplitude of the superconducting band,  $\Delta C/C$  is the relative jump in electronic specific heat at the transition temperature,  $T_c$ , and  $\eta = 2/3$  for *s*-wave superconductors [56]. In the result,  $T_c$ ,  $\Delta C/C$ ,  $\Delta(0)$ , and normal-state tunneling resistance,  $R_n$ , of the S/I/S junction, or in the more general case of S/N/S junction, can be deduced by fitting experimental  $I_c(sf, T)$  datasets to Equation (2), for which the full expression is [51]:

$$I_c(sf, T) = \frac{\pi \cdot \Delta(0) \cdot \tanh\left(\frac{\pi \cdot k_B \cdot T_c}{\Delta(0)} \cdot \sqrt{\eta \cdot \left(\frac{\Delta C}{C}\right) \cdot \left(\frac{T_c}{T} - 1\right)}\right)}{2 \cdot e \cdot R_n} \cdot \tanh\left(\frac{\Delta(0) \cdot \tanh\left(\frac{\pi \cdot k_B \cdot T_c}{\Delta(0)} \cdot \sqrt{\eta \cdot \left(\frac{\Delta C}{C}\right) \cdot \left(\frac{T_c}{T} - 1\right)}\right)}{2 \cdot k_B \cdot T}\right), \quad (4)$$

It should be noted that direct experiments performed by Natterer et al. [59] showed that the superconducting gap does exist in graphene, which is in proximity contact with superconducting electrodes. The gap amplitude,  $\Delta(T)$ , has a characteristic decaying length [59], which is the expected behavior from primary idea of the proximity effect [6]. As a direct consequence, clear physical meaning remains for the relative jump in electronic specific heat at the transition temperature,  $\Delta C/C$ , due to this parameter is an essential thermodynamic consequence for the appearance of the superconducting energy gap,  $\Delta(T)$ . As was shown in another study [51],  $\Delta C/C$  is the fastest decaying parameter of the superconducting state in S/N/S junctions, over the junction length,  $L$ , while  $T_c$  is the most robust one.

In References [51,52], it was shown that S/SLG/S and S/Bi<sub>2</sub>Se<sub>3</sub>/S junctions exhibit two-decoupled band superconducting state. Thus, for the general case of *N*-decoupled bands, the temperature-dependent self-field critical current,  $I_c(sf, T)$ , can be described by the following equation:

$$I_c(sf, T) = \sum_{i=1}^N \frac{\pi \cdot \Delta_i(T)}{2 \cdot e \cdot R_{n,i}} \cdot \theta(T_{c,i} - T) \cdot \tanh\left(\frac{\Delta_i(T)}{2 \cdot k_B \cdot T}\right), \quad (5)$$

where the subscript *i* refers to the *i*-band,  $\theta(x)$  is the Heaviside step function, and each band has its own independent parameters of  $T_{c,i}$ ,  $\Delta C_i/C_i$ ,  $\Delta_i(0)$ , and  $R_{n,i}$ . Equation (5) was also used to analyze experimental  $I_c(sf, T)$  data for several S/DCM/S junctions [60].

Titov and Beenakker [53] proposed that  $I_c(sf, T)$  in S/DCM/S junction at the conditions near the Dirac point can be described by the equation:

$$I_c(sf, T) = 1.33 \cdot \frac{e \cdot \Delta(T)}{\hbar} \cdot \frac{W}{\pi \cdot L}, \quad (6)$$

where  $W$  is the junction width. In this paper, analytical equation for the gap (Equation (3) [57]) is substituted in Equation (6):

$$I_c(sf, T) = 1.33 \cdot \frac{e \cdot \Delta(0) \cdot \tanh\left(\frac{\pi \cdot k_B \cdot T_c}{\Delta(0)} \cdot \sqrt{\eta \cdot \left(\frac{\Delta C}{C}\right) \cdot \left(\frac{T_c}{T} - 1\right)}\right)}{\hbar} \cdot \frac{W}{\pi \cdot L}, \quad (7)$$

with the purpose to deduce  $T_c$ ,  $\Delta C/C$ , and  $\Delta(0)$  values in the S/DCM/S junctions from the fit of experimental  $I_c(sf, T)$  datasets to Equation (7). For a general case of *N*-decoupled bands, temperature-dependent self-field critical current  $I_c(sf, T)$  in S/DCM/S junctions can be described by the following equation:

$$I_c(sf, T) = 1.33 \cdot \frac{e}{\pi \cdot \hbar} \cdot \frac{W}{L} \cdot \sum_{i=1}^N \Delta_i(T) \cdot \theta(T_{c,i} - T), \quad (8)$$

Based on a fact that  $W$  and  $L$  can be measured with very high accuracies, Equation (7) has the minimal ever proposed number of free-fitting parameters (which are  $T_c$ ,  $\Delta C/C$ ,  $\Delta(0)$ ) to fit to the experimental  $I_c(sf, T)$  dataset. However, as we demonstrate below, the ballistic model (Equation (6) [53]) is not the most correct model to describe  $I_c(sf, T)$  in S/DCM/S junctions. It should be noted that Equation (4) utilizes the same minimal set of parameters within the Bardeen-Cooper-Schrieffer (BCS) theory [60],

i.e.,  $T_c$ ,  $\Delta C/C$ ,  $\Delta(0)$ , to describe superconducting state in S/N/S junction and  $R_n$  as a free-fitting parameter to describe the junction.

It should be stressed that a good reason must be presented for requiring a more complex model than is needed to adequately explain the experimental data [61,62].

In the next section, Equations (4), (5), (7), and (8) will be applied to fit experimental  $I_c(sf, T)$  datasets for a variety of S/DCM/S junctions with the purpose to reveal the primary superconducting parameters of these systems and by comparison deduced parameters with weak-coupling s-wave BCS limits we show that the modified Ambegaokar and Baratoff model (Equations (4) and (5)) [51,52] describes the superconducting state in S/DCM/S junctions with higher accuracy.

### 3. Results

#### 3.1. Micrometer-Long Tantalum/Graphene/Tantalum (Ta/G/Ta) Junction

Jang and Kim [63] reported experimental  $I_c(sf, T)$  datasets and fit to KO-1 model (in their Figure 2d [63]) for micrometer long ballistic Ta/G/Ta junctions. The  $I_c(sf, T)$  fit to KO-1 model (Figure 2d [63]) and deduced parameters are in disagreement with experimental values based on  $I_c R_n$  product. In Figure 1, we show  $I_c(sf, T)$  datasets for Device 1 [63] (recorded at gate voltage  $V_g = 10$  V) and fits to single-band ballistic model, Equation (7) (in Figure 1a) and single-band modified AB model Equation (4) (Figure 1b). Device 1 has  $W = 6$   $\mu\text{m}$ ,  $L = 1$   $\mu\text{m}$ , and  $\xi_s = 16$   $\mu\text{m}$  [63]. This means that the ballistic limit of  $L \ll \xi_s$  is satisfied for these junctions.

Results of fits to both models are presented in Table 1.

**Table 1.** Deduced parameters for tantalum/graphene/tantalum (Ta/G/Ta) junction from fit to single-band Titov and Beenakker (TB) [53] and Ambegaokar and Baratoff (AB) [7,8] models.

Parameter	TB Model	AB Model
$T_c$ (K)	$1.052 \pm 0.002$	$1.06 \pm 0.01$
$\Delta C/C$	$17.7 \pm 0.6$	$1.15 \pm 0.07$
$\Delta(0)$ (meV)	$1.03 \pm 0.01$	$0.095 \pm 0.002$
$2\Delta(0)/k_B T_c$	$22.7 \pm 0.3$	$2.1 \pm 0.1$

Deduced parameters from the fit to ballistic model (Equation (7)) in Figure 1a are in remarkable disagreement with any physical-background expectations, i.e., the ratio of  $\frac{2\Delta(0)}{k_B T_c} = 22.7$  (which should be comparable with s-wave BCS weak coupling limit of  $\frac{2\Delta(0)}{k_B T_c} = 3.53$ ) and  $\frac{\Delta C}{C} = 17.7$  (which should be comparable with s-wave BCS weak coupling limit of  $\frac{\Delta C}{C} = 1.43$ ).

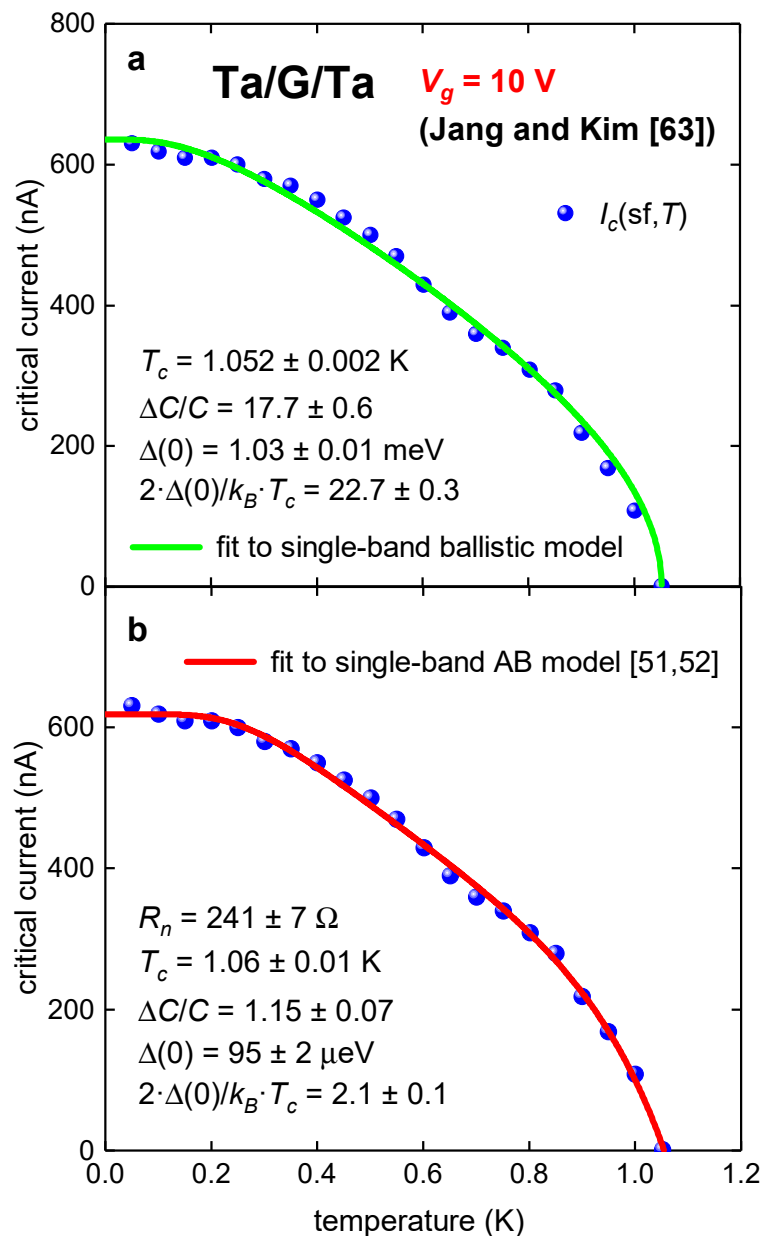
It needs to be noted that the highest experimental value for phonon-mediated superconductors of  $\frac{2\Delta(0)}{k_B T_c} \approx 5$  was measured for lead- and bismuth-based alloys [64,65], and the deduced value by the ballistic model  $\frac{2\Delta(0)}{k_B T_c} \approx 23$  does not have a physical interpretation.

In contrast, the fit to Equation (4) reveals superconducting parameters in expected ranges of  $\frac{2\Delta(0)}{k_B T_c} = 2.1 \pm 0.1$  and  $\frac{\Delta C}{C} = 1.15 \pm 0.07$ , i.e., these parameters are slightly suppressed from s-wave BCS weak-coupling limits as expected [52,60]. It should also be noted that free-fitting parameter  $R_n = 241 \pm 7$   $\Omega$  is in a good agreement with experimental measured value for this junction [63].

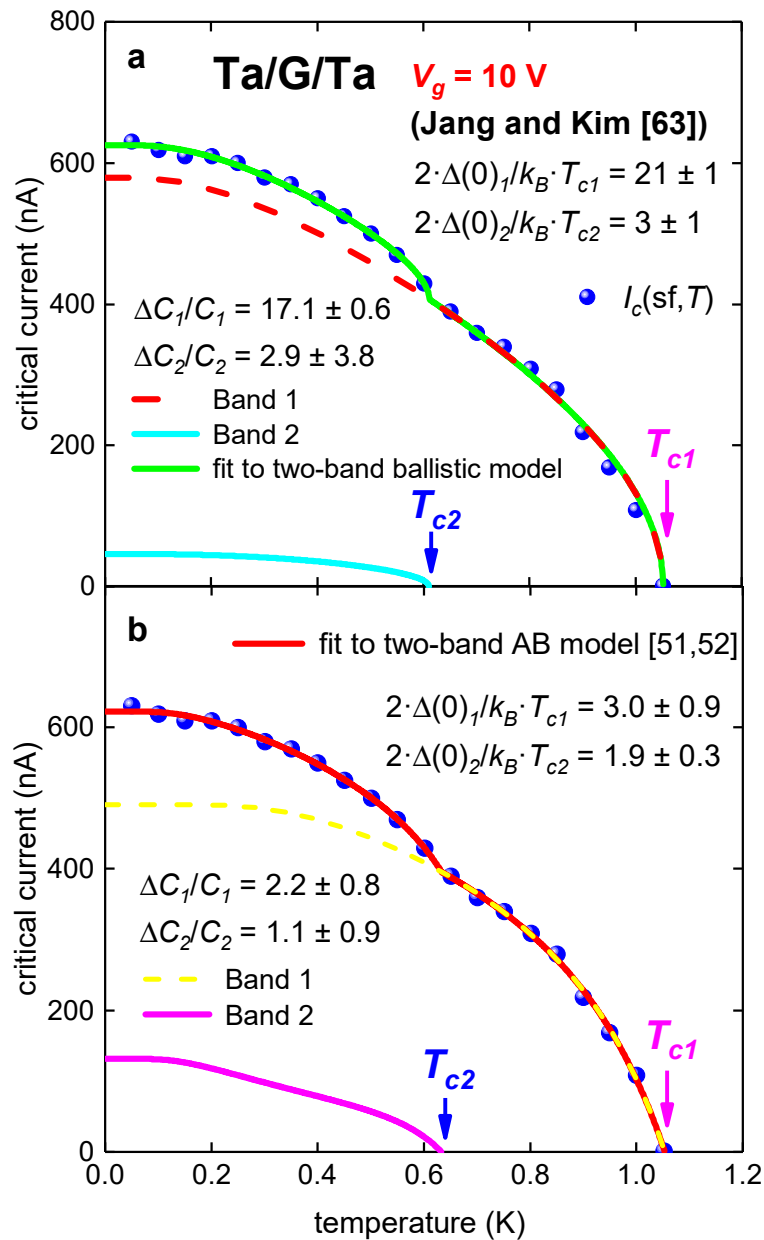
It can be seen (Figure 1), that there is an upturn in experimental  $I_c(sf, T)$  at  $T \sim 0.65$  K, which is a manifestation of the second superconducting band opening in this atomically thin S/N/S junction [51,52]. Thus, the experimental  $I_c(sf, T)$  dataset was fitted to two-band models (Equations (8) and (5)). Results of these fits are shown in Figure 2 and deduced parameters are in Table 2.

**Table 2.** Deduced parameters for tantalum/graphene/tantalum (Ta/G/Ta) junction at  $V_g = 10$  V from fit to two-band Titov and Beenakker (TB) [53] and Ambegaokar and Baratoff (AB) [7,8] models.

Parameter	TB Model	AB Model
$T_{c1}$ (K)	$1.052 \pm 0.001$	$1.053 \pm 0.003$
$T_{c2}$ (K)	$0.61 \pm 0.02$	$0.63 \pm 0.03$
$\Delta C_1/C_1$	$17.1 \pm 0.6$	$2.2 \pm 0.8$
$\Delta C_2/C_2$	$2.9 \pm 3.8$	$1.1 \pm 0.9$
$2 \cdot \Delta_1(0)/k_B \cdot T_{c1}$	$21 \pm 1$	$3.0 \pm 0.9$
$2 \cdot \Delta_2(0)/k_B \cdot T_{c2}$	$3 \pm 1$	$1.9 \pm 0.3$



**Figure 1.** Experimental  $I_c(sf, T)$  for tantalum/graphene/tantalum (Ta/G/Ta) junction (Device 1) at gate voltage of  $V_g = 10$  V [63] and data fits to single-band ballistic model (Equation (7), Panel a) and single-band modified AB model (Equation (4), Panel b) (a) Ballistic model. fit quality is  $R = 0.9948$ ; (b) modified AB model [51,52] fit quality is  $R = 0.9980$ .



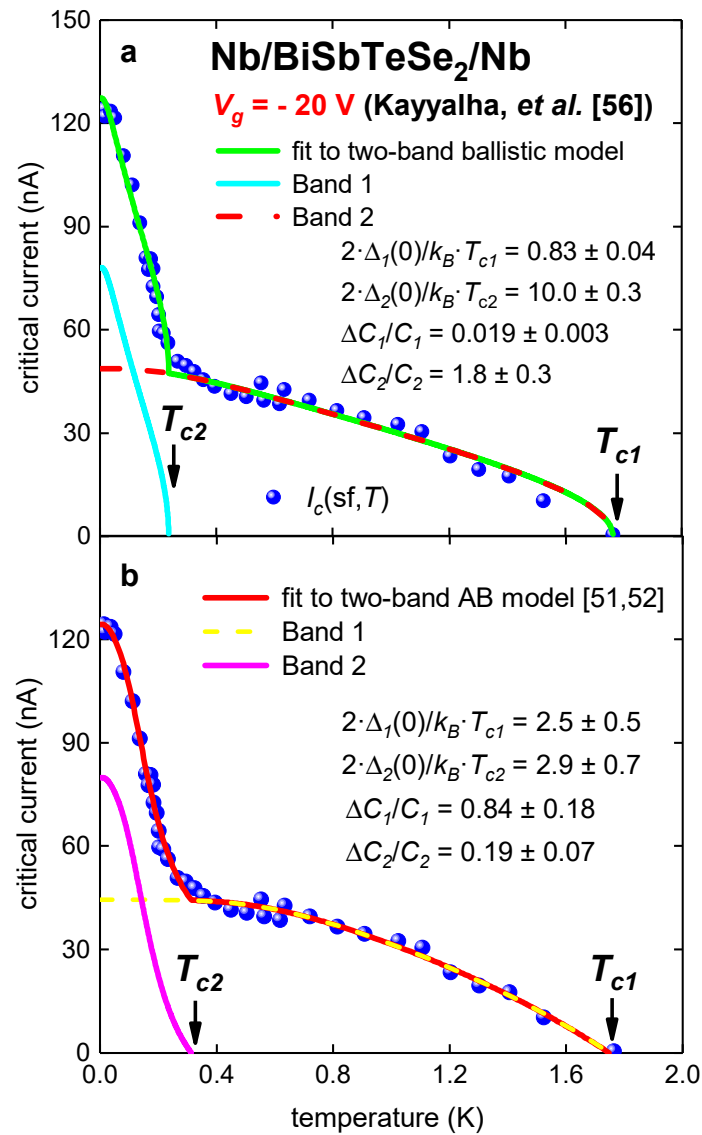
**Figure 2.** Experimental  $I_c(sf, T)$  for Ta/G/Ta junction (Device 1) at gate voltage of  $V_g = 10$  V [63] and data fits to two-band ballistic model (Equation (8), Panel a) and two-band modified AB model (Equation (5), Panel b). (a) Ballistic model, fit quality is  $R = 0.9978$ ; (b) modified AB model [51,52]. Derived parameters:  $R_{n1} = 429 \pm 184 \Omega$ ,  $R_{n2} = 603 \pm 209 \Omega$ , fit quality is  $R = 0.9994$ .

The fit reveals a large disagreement of parameters deduced by ballistic model with expected values within frames for BCS theory. In contrast with this, deduced parameters by modified AB model [51,52] are within weak-coupling limits of BCS. As shown in Reference [51], raw experimental  $I_c(sf, T)$  datasets should be reasonably dense to deduce parameters by AB model with small uncertainties.

### 3.2. Planar Nb/BiSbTeSe<sub>2</sub>-Nanoribbon/Nb Junctions

Kayyalha et al. [56] reported  $I_c(sf, T)$  for five Nb/BiSbTeSe<sub>2</sub>-nanopribbon/Nb junctions at different gate voltage,  $V_g$ . In this paper  $I_c(sf, T)$  datasets for Sample 1 at  $V_g = -20$  V, 0 V and 45 V [56] were analyzed by two-band models (Equations (5) and (8)), because it was already shown in Reference [60] that these junctions exhibit two-band superconducting state. In Figure 3 experimental  $I_c(sf, T)$  dataset [56] and

fits are shown. For this junction,  $L = 40$  nm [56] and  $\xi_s = 640$  nm [56]; thus, the ballistic regime,  $L \ll \xi_s$ , is well satisfied.



**Figure 3.** Experimental  $I_c(sf, T)$  for Nb/BiSbTeSe<sub>2</sub>-nanoribbon/Nb junction (Sample 1 [56]) at gate voltage  $V_g = -20$  V. (a) Ballistic model, fit quality is  $R = 0.990$ ; (b) modified AB model [51,52]. Derived parameters:  $R_{n1} = 6.7 \pm 1.6$  k $\Omega$ ,  $R_{n2} = 0.75 \pm 0.18$  k $\Omega$ , fit quality is  $R = 0.9953$ .

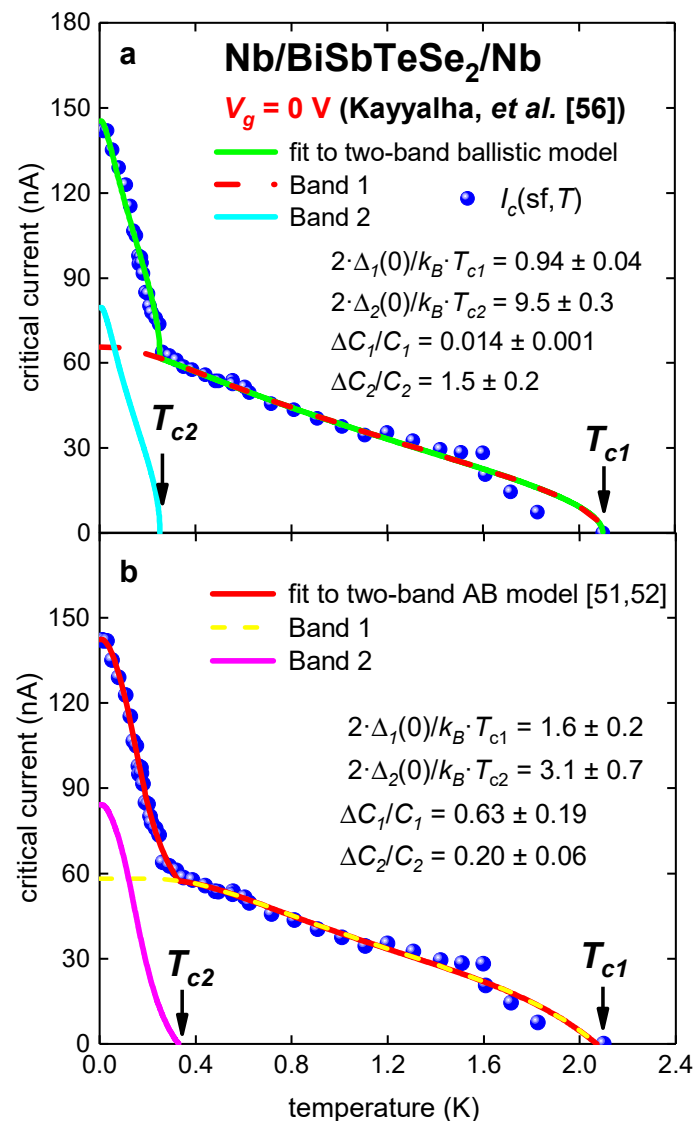
Despite the fact that fits to both models have a similar quality, deduced parameters of the superconducting state (Table 3), i.e.,  $\Delta C_i/C_i$ ,  $\Delta_i(0)$ , and  $\frac{2\Delta_i(0)}{k_B T_{c,i}}$ , for the case of the ballistic models (Figure 3a), similar to the case of Ta/G/Ta junction (Figures 1 and 2), are remarkably different from values expected from BCS theory. Additionally, there are two orders of magnitude difference between deduced  $\Delta C_i/C_i$  for two bands for the same sample, and one order of magnitude for  $\frac{2\Delta_i(0)}{k_B T_{c,i}}$ , which is unavoidable evidence that the ballistic model needs to be reexamined. In contrast with this, the fit to the modified AB model [51] (Figure 3b) reveals deduced parameters, including  $R_{ni}$  values, in the expected ranges. It should be noted that full analysis (within the modified AB model [52]) of  $I_c(sf, T)$  datasets in junctions reported by Kayyalha et al. [56] can be found elsewhere [60].



**Table 3.** Deduced parameters for Nb/BiSbTeSe<sub>2</sub>-nanoribbon/Nb junction (Sample 1 [56]) at  $V_g = -20$  V from fit to two-band Titov and Beenakker (TB) [53] and Ambegaokar and Baratoff (AB) [7,8] models.

Parameter	TB Model	AB Model
$T_{c1}$ (K)	$1.76 \pm 0.01$	$1.74 \pm 0.04$
$T_{c2}$ (K)	$0.236 \pm 0.003$	$0.31 \pm 0.02$
$\Delta C_1/C_1$	$0.019 \pm 0.03$	$0.84 \pm 0.18$
$\Delta C_2/C_2$	$1.8 \pm 0.3$	$0.19 \pm 0.07$
$2 \cdot \Delta_1(0)/k_B \cdot T_{c1}$	$0.83 \pm 0.04$	$2.5 \pm 0.5$
$2 \cdot \Delta_2(0)/k_B \cdot T_{c2}$	$10.0 \pm 0.3$	$2.85 \pm 0.70$

In Figure 4, experimental  $I_c(sf, T)$  dataset [56] and fits to two models for Sample 1 at gate voltage  $V_g = 0$  V also demonstrate that the ballistic model is an inadequate tool to analyze experimental data in S/DCM/S junctions (deduced parameters are given in Table 4).



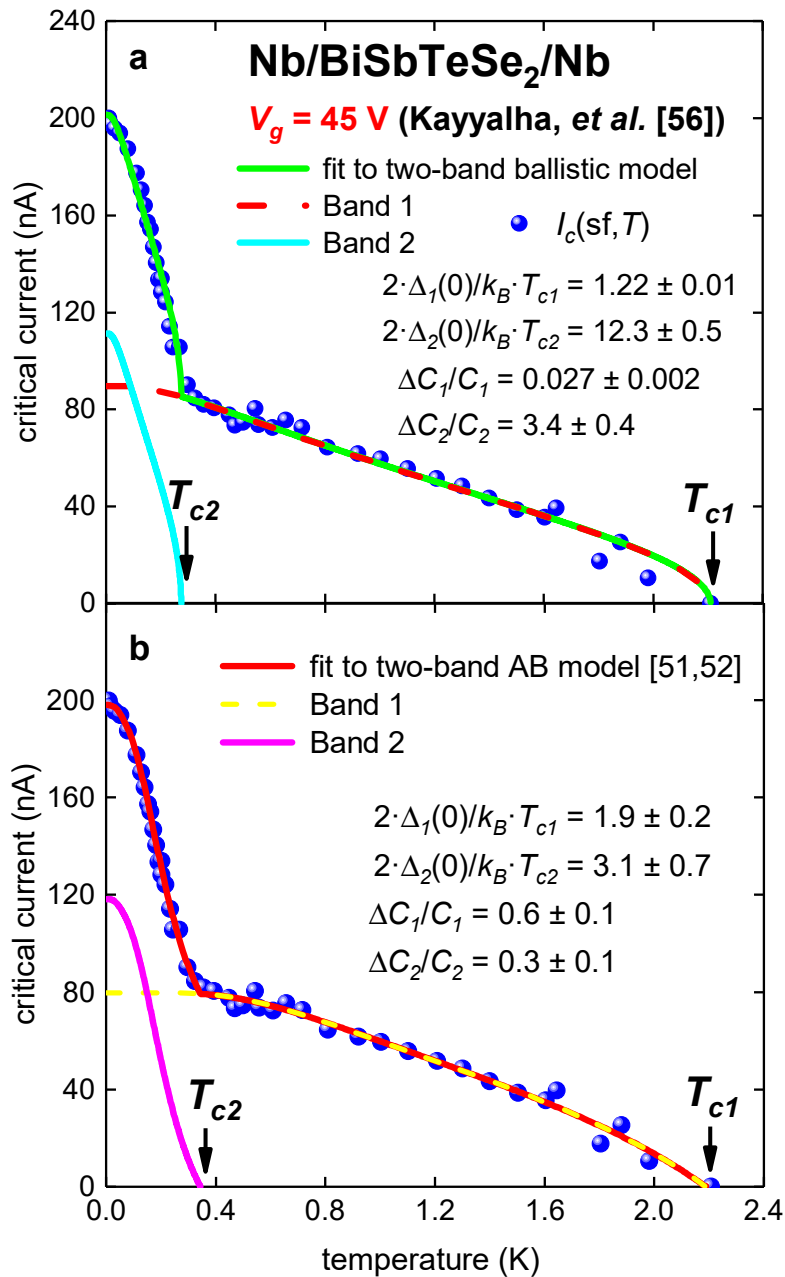
**Figure 4.** Experimental  $I_c(sf, T)$  for Nb/BiSbTeSe<sub>2</sub>-nanoribbon/Nb junction (Sample 1 [56]) at gate voltage  $V_g = 0$  V. (a) Ballistic model, fit quality is  $R = 0.992$ ; (b) modified AB model [51,52]. Derived parameters:  $R_{n1} = 3.9 \pm 0.4$  k $\Omega$ ,  $R_{n2} = 0.81 \pm 0.15$  k $\Omega$ , fit quality is  $R = 0.9965$ .



**Table 4.** Deduced parameters for Nb/BiSbTeSe<sub>2</sub>-nanoribbon/Nb junction (Sample 1 [56]) at  $V_g = 0$  V from fit to two-band Titov and Beenakker (TB) [53] and Ambegaokar and Baratoff (AB) [7,8] models.

Parameter	TB Model	AB Model
$T_{c1}$ (K)	$2.10 \pm 0.01$	$2.07 \pm 0.03$
$T_{c2}$ (K)	$0.252 \pm 0.005$	$0.33 \pm 0.02$
$\Delta C_1/C_1$	$0.014 \pm 0.001$	$0.6 \pm 0.2$
$\Delta C_2/C_2$	$1.5 \pm 0.2$	$0.20 \pm 0.06$
$2 \cdot \Delta_1(0)/k_B \cdot T_{c1}$	$0.94 \pm 0.04$	$1.6 \pm 0.2$
$2 \cdot \Delta_2(0)/k_B \cdot T_{c2}$	$9.5 \pm 0.3$	$3.1 \pm 0.7$

The same conclusion can be made for Sample 1 at  $V_g = 45$  V (Figure 5 and Table 5).



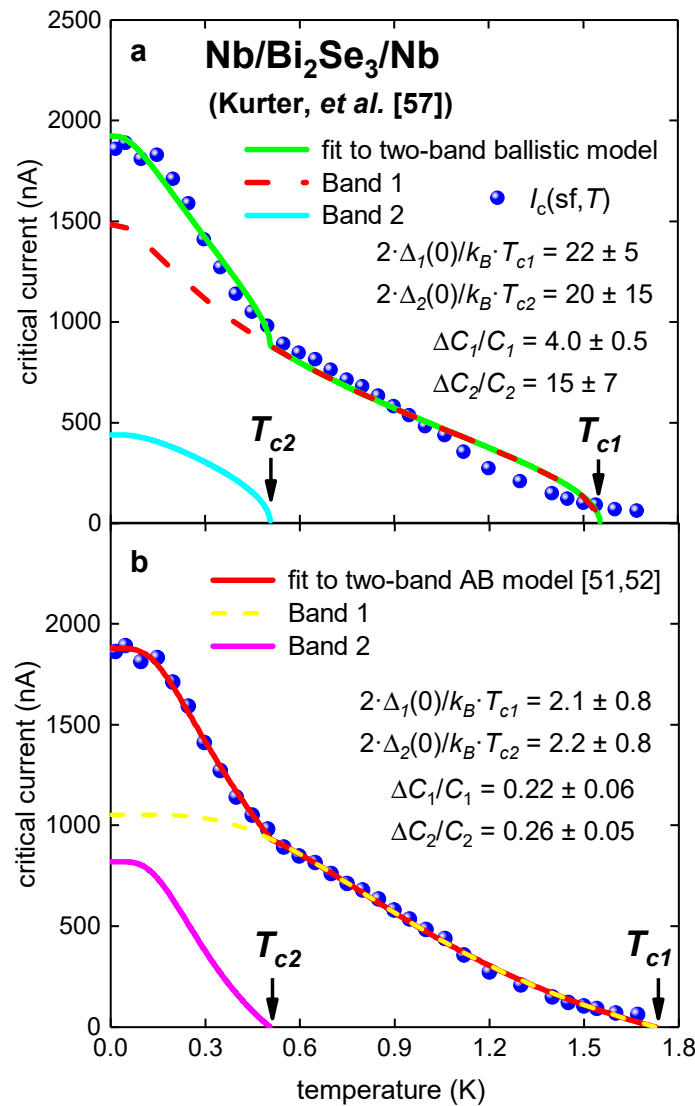
**Figure 5.** Experimental  $I_c(sf, T)$  for Nb/BiSbTeSe<sub>2</sub>-nanoribbon/Nb junction (Sample 1 [56]) at gate voltage  $V_g = 45$  V. (a) Ballistic model, fit quality is  $R = 0.994$ ; (b) modified AB model [51,52]. Derived parameters:  $R_{n1} = 3.5 \pm 0.3$  k $\Omega$ ,  $R_{n2} = 630 \pm 110$   $\Omega$ , fit quality is  $R = 0.998$ .

**Table 5.** Deduced parameters for Nb/BiSbTeSe<sub>2</sub>-nanoribbon/Nb junction (Sample 1 [56]) at  $V_g = 45$  V from fit to two-band Titov and Beenakker (TB) [53] and Ambegaokar and Baratoff (AB) [7,8] models.

Parameter	TB Model	AB Model
$T_{c1}$ (K)	$2.21 \pm 0.01$	$2.19 \pm 0.03$
$T_{c2}$ (K)	$0.274 \pm 0.006$	$0.34 \pm 0.01$
$\Delta C_1/C_1$	$0.027 \pm 0.002$	$0.6 \pm 0.1$
$\Delta C_2/C_2$	$3.4 \pm 0.4$	$0.30 \pm 0.08$
$2 \cdot \Delta_1(0)/k_B \cdot T_{c1}$	$1.22 \pm 0.01$	$1.9 \pm 0.2$
$2 \cdot \Delta_2(0)/k_B \cdot T_{c2}$	$12.3 \pm 0.5$	$3.1 \pm 0.7$

### 3.3. Planar Nb/Bi<sub>2</sub>Se<sub>3</sub>/Nb Junction [56]

In Figure 6, temperature-dependent self-field critical currents,  $I_c(sf, T)$ , in Nb/Bi<sub>2</sub>Se<sub>3</sub>/Nb ( $W = 1000$  nm,  $L = 100$  nm) reported by Kurter et al. [57] is shown. For this junction,  $300 \text{ nm} < \xi_s < 1,000 \text{ nm}$  [57], and thus, the ballistic regime condition,  $L \ll \xi_s$ , is well satisfied.



**Figure 6.** Experimental  $I_c(sf, T)$  for Nb/Bi<sub>2</sub>Se<sub>3</sub>/Nb junction [57]. (a) Ballistic model, fit quality is  $R = 0.994$ ; (b) modified AB model [51,52]. Derived parameters:  $R_{n1} = 240 \pm 100 \Omega$ ,  $R_{n2} = 92 \pm 33 \Omega$ . Fit quality is  $R = 0.9991$ .

There is a large difference between experimental data and the fit to ballistic model (Figure 6 and Table 6). In addition, deduced parameters from the ballistic model fit have no physical interpretation. The fit to the modified AB model reveals parameters in the expected ranges (Figure 6).

**Table 6.** Deduced parameters for Nb/Bi<sub>2</sub>Se<sub>3</sub>/Nb junction [57] from fit to two-band Titov and Beenakker (TB) [53] and Ambegaokar and Baratoff (AB) [7,8] models.

Parameter	TB Model	AB Model
$T_{c1}$ (K)	$1.55 \pm 0.02$	$1.73 \pm 0.05$
$T_{c2}$ (K)	$0.51 \pm 0.03$	$0.51 \pm 0.03$
$\Delta C_1/C_1$	$4.0 \pm 0.5$	$0.22 \pm 0.06$
$\Delta C_2/C_2$	$15 \pm 7$	$0.26 \pm 0.05$
$2 \cdot \Delta_1(0)/k_B \cdot T_{c1}$	$22 \pm 5$	$2.1 \pm 0.8$
$2 \cdot \Delta_2(0)/k_B \cdot T_{c2}$	$15 \pm 7$	$2.2 \pm 0.8$

There is a large difference between experimental data and the fit to ballistic model (Figure 6 and Table 6). In addition, deduced parameters from ballistic model fit have no any physical interpretation. The fit to modified AB model reveals parameters in expected ranges (Figure 6 and Table 6).

#### 4. Discussion

One of the most important questions that can be discussed herein is as follows: what is the origin for such dramatic incapability of ballistic model to analyze the self-field critical currents in S/DCM/S junctions? From the author's point of view, the origin is the primary concept of the KO theory, in that  $I_c(sf, T)$  in the S/N/S junctions is:

$$I_c(sf, T) = \max_{\varphi}(I(\varphi, sf, T)) \quad (9)$$

where  $\varphi$  is the phase difference between two superconducting electrodes of the junction. Despite this assumption is a fundamental conceptual point of the KO theory, there are no physically background or experimental confirmations that this assumption should be a true. In fact, the analysis of experimental data by a model within this assumption (we presented herein) shows that Equation (9) is in remarkably large disagreement with experiment.

One of the simplest ways to show that Equation (9) is incorrect is to note that when the length of the junction,  $L$ , goes to zero, Equation (6) shows:

$$I_c(sf, T) = \lim_{L \rightarrow 0} \left( 1.33 \cdot \frac{e \cdot \Delta(T)}{\hbar} \cdot \frac{W}{\pi \cdot L} \right) \propto \lim_{L \rightarrow 0} \left( \frac{1}{L} \right) \rightarrow \infty. \quad (10)$$

Herein, the simplest available function [53] that was proposed for the S/DCM/S junction in the Equation (9) was chosen as an example. However, other proposed functions for Equation (9) (for which we refer the reader to Reference [12]) have identical unresolved problem, because, as this was shown for about 100 weak-link superconductors [2–5,66], the limit should be (Equation (1)):

$$I_c(sf, T) = \lim_{L \rightarrow 0} \left( 1.33 \cdot \frac{e \cdot \Delta(T)}{\hbar} \cdot \frac{W}{\pi \cdot L} \right) = \frac{\phi_0}{\pi \cdot \mu_0} \cdot \left[ \frac{\ln(1 + \sqrt{2} \cdot \kappa_c(T))}{\lambda_{ab}^3(T)} \cdot \left( \frac{\lambda_c(T)}{b} \cdot \tanh\left(\frac{b}{\lambda_c(T)}\right) \right) + \frac{\ln(1 + \sqrt{2} \cdot \gamma(T) \cdot \kappa_c(T))}{\sqrt{\gamma(T)} \cdot \lambda_{ab}^3(T)} \cdot \left( \frac{\lambda_{ab}(T)}{a} \cdot \tanh\left(\frac{a}{\lambda_{ab}(T)}\right) \right) \right] \cdot (a \cdot b). \quad (11)$$

This means that the primary dissipation mechanism, which governs DC transport current limit in S/N/S, is not yet revealed. However, as we show herein, it is irrelevant to achieving values within the primary concept of KO theory, Equation (9). It should be mentioned that the Density Functional Theory (DFT) calculations [67,68] are currently unexplored powerful techniques, which can be used to reveal dissipation mechanism in S/DCM/S junctions.

## 5. Conclusions

In this paper,  $I_c(sf, T)$  data for S/DCM/S junctions were analyzed by applying two models: the ballistic and the modified Ambegaokar-Baratoff model. It was shown that the ballistic model [10–12,53] cannot describe the self-field critical currents in S/DCM/S junctions. In conclusion, the ballistic model should be reexamined in terms of its applicability to describe dissipation-free self-field transport current in S/DCM/S junctions.

**Funding:** This research was funded by the State Assignment of Minobrnauki of Russia, theme “Pressure” No. AAAA-A18-118020190104-3, and by Act 211 Government of the Russian Federation, contract No. 02.A03.21.0006.

**Conflicts of Interest:** The funders had no role in the design of the study, in the collection, analyses, or interpretation of data, as well as in the writing of the manuscript, or in the decision to publish the results.

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